

A POSTERIORI EIGENVALUE ERROR ESTIMATION FOR THE SCHRÖDINGER OPERATOR WITH THE INVERSE SQUARE POTENTIAL

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ABSTRACT. We develop an *a posteriori* error estimate of hierarchical type for Dirichlet eigenvalue problems of the form $(-\Delta + (c/r)^2)\psi = \lambda\psi$ on bounded domains Ω , where r is the distance to the origin, which is assumed to be in $\bar{\Omega}$. This error estimate is proven to be asymptotically identical to the eigenvalue approximation error on a family of geometrically-graded meshes. Numerical experiments demonstrate this asymptotic exactness in practice.

1. INTRODUCTION

We consider the eigenvalue problem

$$(1) \quad -\Delta\phi + \frac{c^2}{r^2}\phi = \lambda\phi,$$

in a bounded domain $\Omega \subset \mathbb{R}^2$ with the Dirichlet boundary condition $\phi|_{\partial\Omega} = 0$, where $c > 0$ and $r = |x|$ is the distance to the origin, which is assumed to be in $\bar{\Omega}$. This eigenvalue problem is associated with the Schrödinger equation

$$(2) \quad -\Delta u + \frac{c^2}{r^2}u = f,$$

with the Dirichlet boundary condition $u|_{\partial\Omega} = 0$. For simplicity of the theory, we assume that c is constant, and that Ω is a polygon. However, we will also consider non-polygonal domains in the examples as well as in the experiments throughout the paper, to which our results extend. In fact, our analysis can also deal with certain type of non-constant functions c possessing multiple inverse-square singularities, as in [15].

The eigenvalue problem (1) with the inverse-square, or centrifugal, potential $(c/r)^2$ is of importance in quantum mechanics (for example cf.[23, 12, 11, 28]). This potential presents the same “differential order” as the Laplacian near the origin, as is apparent when the Laplacian is expressed in polar coordinates. The strong singularity r^{-2} in the potential generally causes singular behavior (unbounded gradient) in the solution of (2) as well as in some of the eigenfunctions of (1). In addition to the singular potential, the geometry (smoothness) of the domain and boundary conditions may also play a critical role in determining the regularity of the solution. Therefore, new analytical tools, different from techniques for standard elliptic operators with bounded coefficients, are needed to develop well-posedness and regularity results, as well as effective numerical algorithms for (2) and (1). For Schrödinger operators with similar singular potentials, the analysis is generally carried out in Sobolev spaces with special weights, instead of in the usual Sobolev space H^m (for example, see [10, 11, 15, 21, 22] and references therein). In particular, based on the weighted estimates, effective finite element methods associated with a class of graded meshes were proposed in [21] to approximate singular solutions of the Schrödinger equation at the optimal rate. An *a posteriori* error estimate of hierarchical type for these optimal finite element algorithms was developed in [22], and it provides a practical stopping criterion for approximating the solution of (2). The present paper builds on this work, adapting it to eigenvalue problems. We prove, and then numerically demonstrate, that our cheaply-computable error estimate is asymptotically identical to the error in our eigenvalue approximation, independent of singularities present in the eigenfunctions or whether the eigenvalues are degenerate.

Finite element methods for elliptic eigenvalue problems are nearly as old those for the associated boundary value problems, so there is a rich literature, and basic analysis is well-developed, at least for standard

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second-order elliptic operators. We do not attempt a comprehensive overview of the relevant literature, but merely cite two classic references for the basic theory, [3, 25], and mention some three recent papers concerning *a posteriori* error estimates for lower-order methods which might be most readily compared to our own, [7, 24, 6]. In both [24] and [6], eigenvalue error estimates are developed for standard elliptic operators. These are proven to asymptotically exact under certain assumptions on mesh structure and smoothness of the eigenfunctions. Non-self-adjoint problems having real eigenvalues are also considered in [24]. The work [6] employs hierarchical bases for error estimation, in the same manner as we do here, but the effectivity analysis for boundary value problems is quite different from ours, as is the theoretical bridge between boundary value problems and eigenvalue problems—which is done here via a key identity (Lemma 3.1). A certain class of non-linear eigenvalue problems, also relevant in certain quantum physical applications, is considered in [7]. Asymptotic exactness of the eigenvalue errors is not considered in [7], and cannot be achieved for the type of error estimates used, but the important issue of proving convergence of the associated adaptive method is addressed.

The rest of the paper is organized as follows. In Section 2, we define the weighted Sobolev spaces used in the analysis of [21, 22], state key regularity results, and present basic eigenvalue theory for (1). Two examples are presented to provide some intuition about these eigenvalue problems, and one of these is revisited explicitly in the experiments. In Section 3, we first formulate the finite element approximation of the eigenvalue problem on graded meshes (Definition 3.2). Then, using finite element analysis in weighted spaces, we prove the exactness of our *a posteriori* error estimate in Theorem 3.6, our main result. In Section 4, we report numerical tests for different domains with different singular eigenfunctions. These tests confirm our theoretical prediction on the effectivity of the *a posteriori* estimate.

2. BASIC DEFINITIONS AND RESULTS

Throughout, we use the following notation for the L^2 -inner-product and norm,

$$(3) \quad (u, v) = \int_{\Omega} uv \quad , \quad \|u\| = \sqrt{(u, u)} .$$

For multi-indices $\alpha = (\alpha_1, \alpha_2) \in \mathbb{N}_0^2$, we employ the standard conventions $|\alpha| = \alpha_1 + \alpha_2$, and for $v = v(x_1, x_2)$ $\partial^\alpha v = \frac{\partial^{|\alpha|} v}{\partial x_1^{\alpha_1} \partial x_2^{\alpha_2}}$. Let Q consist of the origin and the corners of Ω . These are the points at which one might expect an eigenfunction of (1) to have an unbounded gradient (cf. [1, 2, 13, 14, 27, 17, 18, 16, 20, 8, 4]). For $x \in \overline{\Omega}$, let $\rho(x)$ be the distance between x and Q . We define the following weighted Sobolev spaces and their corresponding norms

$$(4) \quad \mathcal{K}_a^m = \{v \in L^2(\Omega) : \rho^{|\alpha|-a} \partial^\alpha v \in L^2(\Omega) \text{ for all } |\alpha| \leq m\} ,$$

$$(5) \quad |v|_{\mathcal{K}_a^m} = \left(\sum_{|\alpha|=m} \|\rho^{m-a} \partial^\alpha v\|^2 \right)^{1/2} , \quad \|v\|_{\mathcal{K}_a^m}^2 = \left(\sum_{|\alpha| \leq m} |v|_{\mathcal{K}_a^m}^2 \right)^{1/2} .$$

We note that $\mathcal{K}_0^0 = L^2(\Omega)$. Letting

$$(6) \quad \mathcal{H} = \{v \in \mathcal{K}_1^1 : v = 0 \text{ on } \partial\Omega \text{ in the trace sense}\} ,$$

we define the following bilinear form on \mathcal{H} ,

$$(7) \quad B(u, v) = \int_{\Omega} \nabla u \cdot \nabla v + \frac{c^2}{r^2} uv ,$$

and note that it is, in fact, an inner-product. We denote the induced norm by $\|\cdot\|$. It can be shown (cf. [22]) that

Lemma 2.1. *The norms $\|\cdot\|$ and $\|\cdot\|_{\mathcal{K}_1^1}$ are equivalent on \mathcal{H} .*

With these definitions in hand, the variational form of our eigenvalue problem is given by

$$(8) \quad \text{Find } (\lambda, \phi) \in \mathbb{R}_+ \times (\mathcal{H} \setminus \{0\}) \text{ such that } B(\phi, v) = \lambda(\phi, v) \text{ for all } v \in \mathcal{H} .$$

We will refer to a solution (λ, ϕ) of (8) as an eigenpair of B on \mathcal{H} , with eigenvalue λ and eigenfunction ϕ . Before stating a few basic facts about the eigenvalue problem (8), we introduce a related family of boundary value problems

$$(9) \quad \text{Given } f \in L^2(\Omega) \text{ find } u(f) \in \mathcal{H} \text{ such that } B(u(f), v) = (f, v) \text{ for all } v \in \mathcal{H} ,$$

and remark on their well-posedness. Lemma 2.1 leads to the well-posedness of (9) in \mathcal{H} by the Riesz Representation Theorem.

A stronger regularity result is proven for the boundary value problem (9) in [21, Theorem 3.3]:

Theorem 2.2. *There is a constant $\eta > 0$ depending only on Ω and the constant $c \geq 0$ in (7) such that, for any $f \in \mathcal{K}_{a-1}^{m-1}$, where $m \in \mathbb{N}_0$ and $|a| < \eta$, we have $u(f) \in \mathcal{K}_{a+1}^{m+1}$. More specifically, it holds that*

$$(10) \quad \|u(f)\|_{\mathcal{K}_{a+1}^{m+1}} \leq C \|f\|_{\mathcal{K}_{a-1}^{m-1}} ,$$

where C depends on m and a , but not on f .

Let $K : L^2(\Omega) \rightarrow L^2(\Omega)$ be the solution operator associated with (9), i.e. $Kf = u(f)$. Theorem 2.2 implies that K is bounded. The fact that $H_0^1(\Omega)$ is compactly embedded in $L^2(\Omega)$ easily implies that $\mathcal{H} \subset H_0^1(\Omega)$ is compactly embedded in $L^2(\Omega)$, because any bounded sequence in \mathcal{H} is clearly bounded in $H_0^1(\Omega)$ as well. Therefore, K is a self-adjoint compact operator on $L^2(\Omega)$, and we have the following basic results for the eigenvalue problem associated with its inverse – the operator defined by the bilinear form B on \mathcal{H} :

- (1) The eigenvalues of (8) form a sequence of positive numbers with no finite accumulation points. We will assume that they are ordered

$$0 < \lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots .$$

- (2) It is possible to choose a corresponding sequence of eigenfunctions ϕ_n , i.e. $B(\phi_n, v) = \lambda_n(\phi_n, v)$ for all $v \in \mathcal{H}$, such that $(\phi_i, \phi_j) = \delta_{ij}$. Furthermore these eigenfunctions form a Hilbert basis for $L^2(\Omega)$.
- (3) There is a min-max variational characterization of the eigenvalues

$$\lambda_n = \min_{\substack{S \subset \mathcal{H} \\ \dim(S)=n}} \max_{\substack{v \in S \\ v \neq 0}} \frac{\|v\|^2}{\|v\|^2} .$$

Such results are standard for symmetric elliptic problems with L^∞ coefficients.

Given an eigenvalue λ , we denote its invariant subspace by

$$(11) \quad E(\lambda) = \{\psi \in \mathcal{H} : B(\psi, v) = \lambda(\psi, v) \text{ for all } v \in \mathcal{H}\} .$$

In other words, $E(\lambda)$ consists of all the eigenfunctions associated with λ , as well as the zero function. As indicated above by the non-strict inequalities $\lambda_i \leq \lambda_{i+1}$, eigenvalues may be degenerate, having geometric multiplicity $\dim E(\lambda) > 1$.

For bounded domains, it is straight-forward to see that $\mathcal{K}_a^m \subset \mathcal{K}_b^m$ if $b < a$. Let $\eta = \eta(\Omega, c)$ be as in Theorem 2.2, and choose $|a| < \eta$. Let (λ, ϕ) be an eigenpair of (8). Since $\phi \in L^2(\Omega) = \mathcal{K}_0^0 \subset \mathcal{K}_{a-1}^0$, Theorem 2.2 guarantees that $u(\lambda\phi) = \lambda u(\phi) = \phi \in \mathcal{K}_{a+1}^2 \subset \mathcal{K}_{a-1}^2$. By induction, $\phi \in \mathcal{K}_{a+1}^{2n}$ for all $n \in \mathbb{N}$. But $\mathcal{K}_{a+1}^{2n} \subset \mathcal{K}_{a+1}^j$ for any $j \leq 2n$. We therefore have the following corollary.

Corollary 2.3. *There is a constant $\eta > 0$ depending only on Ω and the constant $c \geq 0$ in (7) such that, for any $|a| < \eta$ and any eigenfunction ϕ , it holds that $\phi \in \mathcal{K}_{a+1}^n$ for all $n \geq 0$; more briefly, $\phi \in \mathcal{K}_{a+1}^\infty$.*

We close this section with two examples which help provide some intuition about these eigenvalue problems—particularly the types of singularities which can occur in the eigenfunctions.

Example 2.4. Suppose Ω is the unit disk, $r < 1$. Expressing the eigenvalue problem in polar coordinates and using separation-of-variables, we find the eigenvalues λ_{mn} and corresponding invariant subspaces $E(\lambda_{mn})$ for $n \geq 0$ and $m \geq 1$,

$$\lambda_{mn} = [j_m(\sigma_n)]^2 \quad , \quad E(\lambda_{mn}) = \text{span} \{J_{\sigma_n}(j_m(\sigma_n)r) \cos(n\theta) , J_{\sigma_n}(j_m(\sigma_n)r) \sin(n\theta)\} ,$$

where $j_m(\nu)$ is the m^{th} positive root of the first-kind Bessel function $J_\nu(z)$, and $\sigma_n = \sqrt{n^2 + c^2}$. When $n = 0$ these subspaces are one-dimensional. These formulas hold for $c \geq 0$, but we will primarily be interested in the case $c \in (0, 1)$.

λ	n	m	mult	λ	n	m	mult
9.8696044010893586188	0	1	1	11.394747278578650551	0	1	1
15.920513426475879895	1	1	2	16.823380260414901268	1	1	2
27.181727337203603368	2	1	2	27.799823099432368260	2	1	2
39.478417604357434475	0	2	1	41.856135733780468863	3	1	2
41.354888262245568479	3	1	2	42.644242596364950600	0	2	1

TABLE 1. The smallest eight eigenvalues for the unit disk problem, Example 2.4, listed together with their indices and multiplicities: $c = 1/2$ (left) and $c = 2/3$ (right).

λ	n	m	mult	λ	n	m	mult
9.8696044010893586188	1	1	1	11.394747278578650551	1	1	1
14.681970642123893257	2	1	1	18.278538262077375859	2	1	1
20.190728556426629975	3	1	1	26.374616427163390770	3	1	1
26.374616427163390770	4	1	1	35.642557845428184984	4	1	1
33.217461914268368860	5	1	1	42.644242596364950600	1	2	1
39.478417604357434475	1	2	1	46.052882654426898622	5	1	1
40.706465818200319742	6	1	1	56.113114813020558488	2	2	1
48.831193643619198878	7	1	1	57.582940903291124744	6	1	1

TABLE 2. The smallest eight eigenvalues for the unit sector problem, Example 2.5, listed together with their indices and multiplicities for $c = 0$; $\alpha = 1/2$ (left) and $\alpha = 2/3$ (right).

Since $J_\nu(z) \sim \frac{z^\nu}{2^\nu \Gamma(\nu+1)}$ as $z \rightarrow 0$, eigenfunctions in $E(\lambda_{m0})$ have asymptotic behavior r^c near the origin for this problem, so the gradient will be unbounded at the origin if $c \in (0, 1)$. If $n \geq 1$, the gradient of an eigenfunction in $E(\lambda_{mn})$ vanishes at the origin. Determining the location in the spectrum of all eigenfunctions having a specific regularity is untenable as it would require knowledge of the interlacing of roots of the Bessel functions J_{σ_n} . However, a couple of specific instances will shed light on typical behavior. Table 1 gives the smallest eight eigenvalues (counting multiplicities) for this problem when $c = 1/2$ and $c = 2/3$. These eigenvalues are correct in all digits shown, up to rounding in the last digit.

Example 2.5. Fixing $\alpha \geq 1/2$, suppose Ω is the sector of the unit disk, with $r < 1$ and $0 < \theta < \pi/\alpha$, where θ is the opening angle of the sector. The limiting case $\alpha = 1/2$ represents the unit disk with the positive x -axis removed. As before, we find the eigenvalues λ_{mn} and corresponding invariant subspaces $E(\lambda_{mn})$ for $m, n \geq 1$,

$$\lambda_{mn} = [j_m(\sigma_n)]^2 \quad , \quad E(\lambda_{mn}) = \text{span} \{J_{\sigma_n}(j_m(\sigma_n)r) \sin(n\alpha\theta)\} \quad ,$$

where and $\sigma_n = \sqrt{(n\alpha)^2 + c^2}$. Again, these formulas hold for $c \geq 0$, and the case $c = 0$ (the Laplace eigenvalue problem) illustrates the type of singularities which can occur solely because of re-entrant corners, i.e. $\alpha < 1$. We provide the first eight eigenvalues for $c = 0$, with $\alpha = 1/2$ (slit disk) and $\alpha = 2/3$ (L-shape) in Table 2.

3. DISCRETIZATION AND ERROR ESTIMATION

In this section, we consider the finite element approximation of solutions to the eigenvalue problem (8), with focus on the estimation of error in the computed eigenvalue approximations. Before getting into the details of our finite element discretization, we make a few relevant claims which hold more generally. We restrict our attention to finite dimensional subspaces $V \subset \mathcal{H}$. The natural analogues of (8) and (9) are

$$(12) \quad \text{Find } (\hat{\lambda}, \hat{\phi}) \in \mathbb{R}_+ \times (V \setminus \{0\}) \text{ such that } B(\hat{\phi}, v) = \hat{\lambda}(\hat{\phi}, v) \text{ for all } v \in V \quad ,$$

$$(13) \quad \text{Given } f \in L^2(\Omega) \text{ find } \hat{u}(f) \in V \text{ such that } B(\hat{u}(f), v) = (f, v) \text{ for all } v \in V \quad .$$

As before, we will refer to a solution $(\hat{\lambda}, \hat{\phi})$ of (12) as an eigenpair of B on V , with eigenvalue $\hat{\lambda}$ and eigenfunction $\hat{\phi}$. These discrete problems are well-posed by basic linear algebra and by the coercivity of the

bilinear form on V (Lemma 2.1). More specifically, if $\{v_1, v_2, \dots, v_N\}$ is a basis for V , then (12) is equivalent to the generalized eigenvalue problem

$$(14) \quad A\mathbf{x} = \hat{\lambda}M\mathbf{x} \quad , \quad a_{ij} = B(v_j, v_i) \quad , \quad m_{ij} = (v_j, v_i) \quad ,$$

where the matrices $A = (a_{ij})$ and $M = (m_{ij})$. The following analogues from the continuous eigenvalue problem apply:

- (1) There are precisely $N = \dim(V)$ eigenvalues for the system (14), which we take to be ordered as

$$0 < \hat{\lambda}_1 \leq \hat{\lambda}_2 \leq \dots \leq \hat{\lambda}_N \quad .$$

- (2) It is possible to choose corresponding eigenfunctions $\hat{\phi}_n$, i.e. $B(\hat{\phi}_n, v) = \hat{\lambda}_n(\hat{\phi}_n, v)$ for all $v \in V$, such that $(\hat{\phi}_i, \hat{\phi}_j) = \delta_{ij}$. These eigenfunctions clearly form a Hilbert basis for V .
- (3) There is a min-max variational characterization of the eigenvalues

$$\hat{\lambda}_n = \min_{\substack{S \subset V \\ \dim(S)=n}} \max_{\substack{v \in S \\ v \neq 0}} \frac{\|v\|^2}{\|v\|^2} = \min_{\substack{S \subset \mathbb{R}^n \\ \dim(S)=n}} \max_{\substack{\mathbf{x} \in S \\ \mathbf{x} \neq 0}} \frac{\mathbf{x}^t A \mathbf{x}}{\mathbf{x}^t M \mathbf{x}} \quad .$$

This characterization implies that $\hat{\lambda}_n \geq \lambda_n$ for $n \leq N$.

Lemma 3.1. *Suppose that (λ, ϕ) is an eigenpair for B on \mathcal{H} and $(\hat{\lambda}, \hat{\phi})$ is an eigenpair for B on V , with $\|\phi\| = \|\hat{\phi}\| = 1$. Let $\tilde{\phi} = u(\hat{\lambda}\hat{\phi})$. It holds that*

$$(15) \quad \hat{\lambda} - \lambda = \|\tilde{\phi} - \hat{\phi}\|^2 + (\lambda - \hat{\lambda}) \left[(\tilde{\phi} - \hat{\phi}, \phi) + (\hat{\phi}, \phi - \hat{\phi}) \right] + \hat{\lambda}(\tilde{\phi} - \hat{\phi}, \phi - \hat{\phi}) \quad .$$

Proof. Using the fact that $\|\phi\| = \|\hat{\phi}\| = 1$, we first have the following identities,

$$\begin{aligned} \|\phi - \hat{\phi}\|^2 &= B(\phi - \hat{\phi}, \phi - \hat{\phi}) = \lambda + \hat{\lambda} - 2\lambda(\phi, \hat{\phi}) \quad , \\ \lambda\|\phi - \hat{\phi}\|^2 &= \lambda(\phi - \hat{\phi}, \phi - \hat{\phi}) = \lambda + \lambda - 2\lambda(\phi, \hat{\phi}) \quad . \end{aligned}$$

Subtracting these identities, we obtain the well-known error formula,

$$(16) \quad \hat{\lambda} - \lambda = \|\phi - \hat{\phi}\|^2 - \lambda\|\phi - \hat{\phi}\|^2 \quad .$$

Using the fact that $B(\tilde{\phi}, v) = \hat{\lambda}(\hat{\phi}, v)$ for all $v \in \mathcal{H}$, we further manipulate (16),

$$\begin{aligned} \hat{\lambda} - \lambda &= B(\phi - \hat{\phi}, \phi - \hat{\phi}) - \lambda(\phi - \hat{\phi}, \phi - \hat{\phi}) \\ &= \lambda(\phi, \phi - \hat{\phi}) - B(\hat{\phi}, \phi - \hat{\phi}) - \lambda(\phi, \phi - \hat{\phi}) + \lambda(\hat{\phi}, \phi - \hat{\phi}) \\ &= \hat{\lambda}(\hat{\phi}, \phi - \hat{\phi}) - B(\hat{\phi}, \phi - \hat{\phi}) + (\lambda - \hat{\lambda})(\hat{\phi}, \phi - \hat{\phi}) \\ &= B(\tilde{\phi} - \hat{\phi}, \phi - \hat{\phi}) + (\lambda - \hat{\lambda})(\hat{\phi}, \phi - \hat{\phi}) \\ &= \|\tilde{\phi} - \hat{\phi}\|^2 + B(\tilde{\phi} - \hat{\phi}, \phi - \tilde{\phi}) + (\lambda - \hat{\lambda})(\hat{\phi}, \phi - \hat{\phi}) \\ &= \|\tilde{\phi} - \hat{\phi}\|^2 + \lambda(\tilde{\phi} - \hat{\phi}, \phi) - \hat{\lambda}(\tilde{\phi} - \hat{\phi}, \hat{\phi}) + (\lambda - \hat{\lambda})(\hat{\phi}, \phi - \hat{\phi}) \quad , \end{aligned}$$

from which (15) follows directly. □

Our computed estimate of $\lambda - \hat{\lambda}$ will be based on approximating $\|\tilde{\phi} - \hat{\phi}\|^2$, treating the rest of the bound in (15) as higher-order terms. To make this more precise, we now shift to definitions of our finite element spaces, and a few key results.

Given a triangulation \mathcal{T} of Ω , let \mathcal{V} be the vertex set (the vertices of all triangles), which we assume includes all singular points Q . We define the two spaces

$$(17) \quad V = V(\mathcal{T}) = \{\mathcal{H} \cap C(\overline{\Omega}) : v|_T \in \mathbb{P}_1, \forall T \in \mathcal{T}\},$$

$$(18) \quad W = W(\mathcal{T}) = \{\mathcal{H} \cap C(\overline{\Omega}) : v|_T \in \mathbb{P}_2, \forall T \in \mathcal{T} \text{ and } v(z) = 0, \forall z \in \mathcal{V}\}.$$

The space \mathbb{P}_k consists of polynomials of total degree k or less. We note that it is necessary that $v(\mathbf{0}) = 0$ for $v \in V$. We will approximate the solution of (8) in the space V and assess the error of this approximation in the space W .

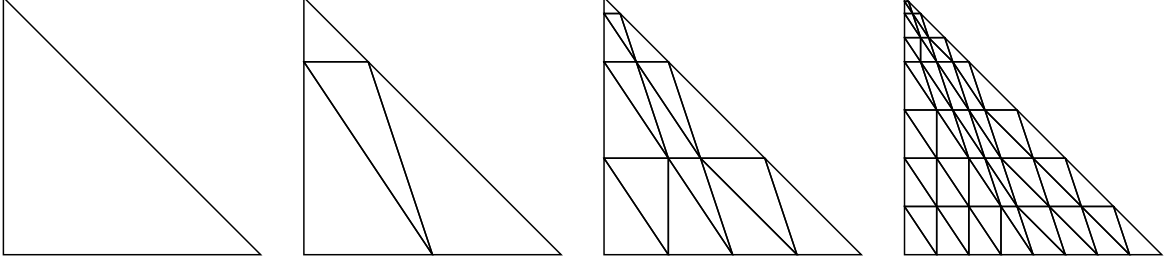


FIGURE 1. A triangle and three consecutive κ -refinements towards the top vertex, $\kappa = 1/4$.

Definition 3.2 (Graded Triangulations). Let \mathcal{T} be a triangulation of Ω whose vertices include Q , such that no triangle in \mathcal{T} has more than one of its vertices in Q . For $\kappa \in (0, 1/2]$, a κ refinement of \mathcal{T} , denoted by $\kappa(\mathcal{T})$, is obtained by dividing each edge AB of \mathcal{T} in two parts as follows:

- If neither A nor B is in Q , then we divide AB into two equal parts.
- Otherwise, if say A is in Q , we divide AB into AC and CB such that $|AC| = \kappa|AB|$.

This will divide each triangle of \mathcal{T} into four triangles. Figure 1 shows a triangle having a singular vertex (the vertex on the top), together with three subsequent κ -refinements, with $\kappa = 1/4$. Given an initial triangulation \mathcal{T}_0 , the associated family of graded triangulations $\{\mathcal{T}_n : n \geq 0\}$ is defined recursively, $\mathcal{T}_{n+1} = \kappa(\mathcal{T}_n)$.

Remark 3.3. Although it may be useful in practice to have a different grading ratio κ_q for each $q \in Q$, which is not difficult to implement, we do not pursue that theoretical generality here.

Given a family $\{\mathcal{T}_n\}$ of κ -refined triangulations, we set $V_n = V(\mathcal{T}_n)$ and $W_n = W(\mathcal{T}_n)$, and $u_n(f) \in V_n$ is the solution of (13) on V_n . We see that $\dim(V_n) \sim \dim(W_n) \sim 4^n$, because each refinement increases the number of triangles by precisely a factor of 4. We also define $\varepsilon_n(f) \in W_n$ by

$$(19) \quad B(\varepsilon_n(f), v) = (f, v) - B(u_n(f), v) \text{ for all } v \in W_n .$$

We collect two key results from [21, 22].

Theorem 3.4. Let η be as in Theorem 2.2, and for $0 < a < \min(\eta, 1)$ choose $\kappa = 2^{-1/a}$. There is a constant C which is independent of f and n , such that

$$(20) \quad \|u(f) - u_n(f)\| \leq C2^{-n} \|f\|_{\mathcal{K}_{a-1}^0}$$

$$(21) \quad \|u(f) - u_n(f) - \varepsilon_n(f)\| \leq C2^{-\sigma n} \|f\|_{\mathcal{K}_{\xi-1}^1}$$

where the related numbers $\sigma > 1$ and $a < \xi < \min(2a, \eta, 1)$ are also independent of f and n .

Although we generally think of f as remaining fixed, these results allow for f to vary with n , and we exploit this fact below. In what follows, we let $\{(\lambda_{k,n}, \phi_{k,n}) : 1 \leq k \leq N = \dim(V_n)\}$, be eigenpairs for B on V_n with $(\phi_{i,n}, \phi_{j,n}) = \delta_{ij}$, as discussed at the beginning of this section; therefore, $\lambda_{k,n} = \hat{\lambda}_k$, $\phi_{k,n} = \hat{\phi}_k$ and $V = V_n$, for example.

Corollary 3.5. Setting $\psi_{k,n} = u(\lambda_{k,n} \phi_{k,n})$ and $\varepsilon_{k,n} = \varepsilon_n(\lambda_{k,n} \phi_{k,n})$, and employing the assumptions of Theorem 3.4, we have

$$(22) \quad \|\psi_{k,n} - \phi_{k,n}\| \leq C\lambda_{k,n} 2^{-n} ,$$

$$(23) \quad \|\psi_{k,n} - \phi_{k,n} - \varepsilon_{k,n}\| \leq C\lambda_{k,n}^{3/2} 2^{-\sigma n} ,$$

where C is independent of k and n .

Proof. Putting (20) in this context, we have

$$\|\psi_{k,n} - \phi_{k,n}\| \leq C2^{-n} \|\lambda_{k,n} \phi_{k,n}\|_{\mathcal{K}_{a-1}^0} \leq C2^{-n} \lambda_{k,n} \|\phi_{k,n}\|_{L^2(\Omega)} \leq C\lambda_{k,n} 2^{-n} .$$

Similarly,

$$\|\psi_{k,n} - \phi_{k,n} - \varepsilon_{k,n}\| \leq C2^{-\sigma n} \|\lambda_{k,n} \phi_{k,n}\|_{\mathcal{K}_{\xi-1}^1} \leq C2^{-\sigma n} \lambda_{k,n} \|\phi_{k,n}\| \leq C\lambda_{k,n}^{3/2} 2^{-\sigma n} .$$

This completes the proof. \square

We emphasize that the computation of $\varepsilon_{k,n}$ involves solving the problem

$$(24) \quad B(\varepsilon_{k,n}, v) = \lambda_{k,n}(\phi_{k,n}, v) - B(\phi_{k,n}, v) \text{ for all } v \in W_n .$$

Using the standard bases for V_n and W_n , it is shown in [22, Theorem 3.6] that the condition number of the matrix associated with (24) is well-conditioned independent of n . In fact, it is spectrally equivalent to its own diagonal. This makes (24) very inexpensive to solve, particularly when compared with computing solutions to the eigenvalue problem $A\mathbf{x} = \hat{\lambda}M\mathbf{x}$ on V_n , where the stiffness matrix A is known to have a condition number which grows like 4^n .

Suppose we fix k and consider the sequence of discrete eigenpairs $\{(\lambda_{k,n}, \phi_{k,n}) : n \geq 0\}$ for the Schrödinger operator with meshes appropriately graded near the set Q (see Theorem 3.4). We do not rehearse standard finite element convergence theory for eigenvalue problems (cf. [9, Section 3.3], [3]), but by the approximation property given in Theorem 3.4, the following results for our eigenvalue problem can be derived in a similar fashion:

- The approximate eigenvalues $\lambda_{k,n}$ converge (down) to λ_k quadratically,

$$(25) \quad \lambda_k \leq \lambda_{k,n} \text{ with } \lambda_{k,n} - \lambda_k = \mathcal{O}(4^{-n}) .$$

- The distance between $\phi_{k,n}$ and the invariant subspace $E(\lambda_k)$ associated with λ_k decreases linearly in the energy norm,

$$(26) \quad \min_{v \in E(\lambda_k)} \|v - \phi_{k,n}\| = \mathcal{O}(2^{-n}) .$$

We emphasize that λ_k may be a degenerate eigenvalue (repeated in the sequence of eigenvalues), so $E(\lambda_k)$ may have dimension greater than one. The analogous statements on V_n hold as well. In light of this, it does not necessarily make sense to say that $\{\phi_{k,n} : n \geq 0\}$ converges, even up to sign. Nevertheless, we do have convergence in the sense of (26), and we refer to this (loosely) as eigenvector “convergence”. We also remark that, although the eigenfunction $v \in E(\lambda_k)$ which is nearest to $\phi_{k,n}$ may not be of unit length, but we do not lose (26) if we add this restriction.

In practice, the eigenvalue convergence is precisely quadratic, and the eigenvector “convergence” is precisely linear on these properly graded meshes.

We now reconsider the various terms in the error identity (15),

$$\begin{aligned} \lambda_{k,n} - \lambda_k &= \|\psi_{k,n} - \phi_{k,n}\|^2 + (\lambda_{k,n} - \lambda_k) [(\psi_{k,n} - \phi_{k,n}, v_{k,n}) + (\phi_{k,n}, v_{k,n} - \phi_{k,n})] \\ &\quad + \lambda_{k,n}(\psi_{k,n} - \phi_{k,n}, v_{k,n} - \phi_{k,n}) . \end{aligned}$$

Here we have taken $v_{k,n} = \operatorname{argmin}\{\|v - \phi_{k,n}\| : v \in E(\lambda_k), \|v\| = 1\}$.

- We take the simple bound $\|v_{k,n} - \phi_{k,n}\| \leq C\|v_{k,n} - \phi_{k,n}\| = \mathcal{O}(2^{-n})$.
- Using a duality argument (L^2 -lifting, or Nitsche’s trick), we see that $\|\psi_{k,n} - \phi_{k,n}\| = \mathcal{O}(4^{-n})$.
- Finally, we note that

$$\begin{aligned} \|\psi_{k,n} - \phi_{k,n}\|^2 - \|\varepsilon_{k,n}\|^2 &= (\|\psi_{k,n} - \phi_{k,n}\| - \|\varepsilon_{k,n}\|)(\|\psi_{k,n} - \phi_{k,n}\| + \|\varepsilon_{k,n}\|) \\ &\leq (\|\psi_{k,n} - \phi_{k,n} - \varepsilon_{k,n}\|)(\|\psi_{k,n} - \phi_{k,n}\| + \|\varepsilon_{k,n}\|) = \mathcal{O}(2^{-(1+\sigma)n}) . \end{aligned}$$

Combining these pieces, we arrive at our key eigenvalue error theorem.

Theorem 3.6. *Under the assumptions of Theorem 3.4, it holds that*

$$(27) \quad \lambda_{k,n} - \lambda_k = \|\varepsilon_{k,n}\|^2 + \mathcal{O}(4^{-\tau n}) ,$$

for some constant $\tau > 1$. The hidden constant in $\mathcal{O}(4^{-\tau n})$ depends on λ_k .

n	N	$\lambda_{1,n} - \lambda_1$	$\ \varepsilon_{1,n}\ ^2$	EFF	$\lambda_{6,n} - \lambda_6$	$\ \varepsilon_{6,n}\ ^2$	EFF
0	48	9.467 ₋₁	8.284 ₋₁	1.142			
1	224	2.429 ₋₁	1.907 ₋₁	1.273	3.892 ₊₀	3.395 ₊₀	1.146
2	961	5.690 ₋₂	4.594 ₋₂	1.238	9.629 ₋₁	8.450 ₋₁	1.139
3	3968	1.631 ₋₂	1.413 ₋₂	1.154	2.493 ₋₁	2.284 ₋₁	1.091
4	16129	3.957 ₋₃	3.605 ₋₃	1.098	6.182 ₋₂	5.898 ₋₂	1.048
5	65025	1.026 ₋₃	9.527 ₋₄	1.076	1.560 ₋₂	1.510 ₋₂	1.033
6	261121	2.637 ₋₄	2.469 ₋₄	1.068	3.929 ₋₃	3.844 ₋₃	1.022

TABLE 3. Data for the Unit Disk problem, corresponding to approximations of λ_1 and λ_6 on graded meshes with $\kappa = 0.2$, $c = 1/2$.

4. NUMERICAL EXPERIMENTS

In this section we report the outcome of several numerical experiments, to demonstrate how well the theory of previous sections—particularly Theorem 3.6—are realized in practice. The data of interest are the eigenvalue errors $\lambda_{k,n} - \lambda_k$, their computed estimates $\|\varepsilon_{k,n}\|^2$, and the associated effectivities

$$EFF = \frac{\lambda_{k,n} - \lambda_k}{\|\varepsilon_{k,n}\|^2}.$$

The software package PLTMG [5] was used for these experiments, with suitable modifications for employing hierarchical basis error estimation and graded mesh refinement, with ARPACK [19] in shift-and-invert mode as the algebraic eigenvalue solver. In order to reduce the width of tables of numerical data, we use the following abbreviation of scientific notation, $a \times 10^m \leftrightarrow a_m$. For example,

$$1.949 \times 10^{-4} \leftrightarrow 1.949_{-4}.$$

We first revisit the unit disk problem of Example 2.4, considering case $c = 1/2$, for which we know that the eigenfunctions associated with $\lambda_1 \approx 9.86960440$ and $\lambda_6 \approx 39.4784176$ have an $r^{1/2}$ singularity (see Table 1). The grading ratio $\kappa = 0.2$ was used for refinement. Note that by Theorem 3.4, the upper bound of the grading parameter κ near the origin is $2^{-1/(1/2)} = 0.25$ to achieve the optimal convergence rate. Therefore, we have chosen an appropriate grading ratio here. The data for these experiments are in Table 3. The eigenvalue convergence is seen to be quadratic, i.e. linear in $N = \dim(V_n)$, and the effectivities are very close 1. The top row of data is absent for λ_6 because, on this coarse mesh, the approximate eigenvalue 33.2876671 was actually (slightly) nearer to $\lambda_4 = \lambda_5 \approx 27.1817273$ than to λ_6 . The effectivity of the error estimate when this was taken into account was 1.010.

We now consider the degenerate eigenvalue $\lambda = \lambda_2 = \lambda_3 \approx 15.9205134$. The corresponding invariant subspace (eigenspace) is spanned by

$$\phi_2 = J_\sigma(\sqrt{\lambda} r) \cos(\theta) \quad , \quad \phi_3 = J_\sigma(\sqrt{\lambda} r) \sin(\theta) \quad \text{where} \quad \sigma = \frac{\sqrt{5}}{2} \approx 1.11803.$$

Of course the ordering of ϕ_2 and ϕ_3 is arbitrary, as is that particular choice of basis for this invariant subspace. These functions are smooth enough to be optimally approximated on a sequence of uniformly refined meshes, $\kappa = 0.5$; for comparison, grading ratio $\kappa = 0.4$ and $\kappa = 0.2$ were used as well. On each mesh, two approximate eigenvalues and eigenvectors were computed and error estimates for both were computed. The results indicate that it really is irrelevant which of the approximate eigenpairs is used to estimate error in the eigenvalue approximation, as indicated by our theory. Since the code (PLTMG+ARPACK) assigns an order to the approximate eigenpairs, we employ this order as well, $(\lambda_{2,n}, \phi_{2,n})$, $(\lambda_{3,n}, \phi_{3,n})$. The computed eigenvalues $\lambda_{2,n}$ and $\lambda_{3,n}$ agreed with each other to far more digits than they agreed with $\lambda_2 = \lambda_3$, so the reported errors are identical. It is only the error estimates, and hence effectivities, which are slightly different. In terms of the grading, all three grading choices gave optimal order convergence, as the theory predicts, with uniform refinement ($\kappa = 0.5$) yielding the smallest errors and $\kappa = 0.2$ yielding the largest errors. In terms of effectivities, uniform refinement was the worst, followed in order by $\kappa = 0.4$ and $\kappa = 0.2$, though all were close to 1. To save space, only the data for $\kappa = 0.5$ and $\kappa = 0.2$ are reported in Table 4. In order to

n	N	$\lambda_{2,n} - \lambda_2$	$\ \varepsilon_{2,n}\ ^2$	EFF	$\lambda_{3,n} - \lambda_3$	$\ \varepsilon_{3,n}\ ^2$	EFF
0	48	1.019 ₊₀	9.181 ₋₁	1.110	1.019 ₊₀	9.181 ₋₁	1.110
1	224	2.544 ₋₁	2.233 ₋₁	1.139	2.544 ₋₁	2.197 ₋₁	1.158
2	961	6.371 ₋₂	5.498 ₋₂	1.158	6.371 ₋₂	5.588 ₋₂	1.140
3	3968	1.596 ₋₂	1.424 ₋₂	1.121	1.596 ₋₂	1.413 ₋₂	1.129
4	16129	3.997 ₋₃	3.562 ₋₃	1.122	3.997 ₋₃	3.576 ₋₃	1.118
5	65025	1.001 ₋₃	8.987 ₋₄	1.114	1.001 ₋₃	8.966 ₋₄	1.116
6	261121	2.506 ₋₄	2.250 ₋₄	1.113	2.506 ₋₄	2.250 ₋₄	1.113

n	N	$\lambda_{2,n} - \lambda_2$	$\ \varepsilon_{2,n}\ ^2$	EFF	$\lambda_{3,n} - \lambda_3$	$\ \varepsilon_{3,n}\ ^2$	EFF
0	48	2.371 ₊₀	2.361 ₊₀	1.004	2.371 ₊₀	2.361 ₊₀	1.004
1	224	5.769 ₋₁	5.170 ₋₁	1.116	5.769 ₋₁	5.272 ₋₁	1.093
2	961	1.433 ₋₁	1.299 ₋₁	1.103	1.433 ₋₁	1.278 ₋₁	1.122
3	3968	3.576 ₋₂	3.311 ₋₂	1.080	3.576 ₋₂	3.282 ₋₂	1.090
4	16129	8.938 ₋₃	8.406 ₋₃	1.063	8.938 ₋₃	8.374 ₋₃	1.067
5	65025	2.234 ₋₃	2.117 ₋₃	1.055	2.234 ₋₃	2.114 ₋₃	1.057
6	261121	5.586 ₋₄	5.312 ₋₄	1.052	5.586 ₋₄	5.308 ₋₄	1.052

TABLE 4. Data for the Unit Disk problem, corresponding to approximations of $\lambda_2 = \lambda_3$. Uniform refinement (top) and $\kappa = 0.2$ graded refinement (bottom), $c = 1/2$.

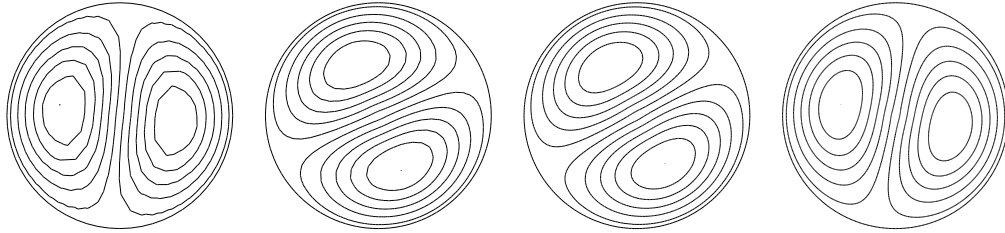


FIGURE 2. Contour plots of $\phi_{2,n}$ for $n = 1, 2, 3, 4$ for the Unit Disk problem, with $\kappa = 0.4$.

demonstrate the “drift” in approximate eigenfunctions associated with degenerate eigenvalues, we provide a sequence of contour plots for $\phi_{2,n}$, $n = 1, 2, 3, 4$ in Figure 2. The contour plots of $\phi_{3,n}$ are essentially obtained by rotating the given plots by 90 degrees. These plots illustrate the assertion in Section 3 that the sequence $\{\phi_{k,n}\}$ may not converge, though the terms are getting successively closer to $E(\lambda_k)$.

Finally, we turn to the L-shape domain $\Omega = (-1, 3)^2 \setminus [1, 3)^2$. We consider the case $c = 1/2$, for which there will be eigenfunctions having an $r^{1/2}$ -singularity at the origin and an $r^{2/3}$ -singularity at the point $(1, 1)$. We use the grading ratios $\kappa = 0.2$ for triangles touching the origin, and $\kappa = 0.3$ for triangles touching $(1, 1)$. Table 5 contains our approximations and error estimates for the first four eigenvalues. Contour plots of the first four eigenfunctions are given in Figure 3. As an interesting comparison, we also consider the case $c = 0$, for which no singularity is present at the origin, and grading is only needed near the point $(1, 1)$. The eigenvalues in this case have been obtained elsewhere to very high accuracy [26] using a computational very well-suited to the Laplacian, and we report their values here, rescaling them by a factor of four due to the fact that our domain has four times the area of theirs:

$$(28) \quad \lambda_1 \approx 2.4099310 \quad , \quad \lambda_2 \approx 3.7993130 \quad , \quad \lambda_3 = \frac{\pi^2}{2} \approx 4.9348022 \quad , \quad \lambda_4 \approx 7.3803703 .$$

The second eigenfunctions for $c = 1/2$ and $c = 0$ are not linearly dependent, nor are fourth eigenfunctions for both $c = 1/2$ and $c = 0$. Their contour plots are merely very similar, though not identical.

n	N	$\lambda_{1,n}$	$\ \varepsilon_{2,n}\ ^2$	$\lambda_{2,n}$	$\ \varepsilon_{2,n}\ ^2$	$\lambda_{3,n}$	$\ \varepsilon_{3,n}\ ^2$	$\lambda_{4,n}$	$\ \varepsilon_{4,n}\ ^2$
0	16	4.416	1.326 ₊₀	5.091	1.163 ₊₀	8.044	2.135 ₊₀	11.00	3.150 ₊₀
1	80	3.491	3.064 ₋₁	4.212	2.682 ₋₁	6.582	5.633 ₋₁	8.488	7.348 ₋₁
2	353	3.251	7.477 ₋₂	4.005	6.746 ₋₂	6.128	1.387 ₋₁	7.878	1.930 ₋₁
3	1472	3.194	2.126 ₋₂	3.953	1.730 ₋₂	6.024	4.031 ₋₂	7.723	5.063 ₋₂
4	6017	3.177	5.588 ₋₃	3.940	4.361 ₋₃	5.992	1.055 ₋₂	7.684	1.283 ₋₂
5	24321	3.173	1.488 ₋₃	3.937	1.094 ₋₃	5.984	2.813 ₋₃	7.674	3.242 ₋₃
6	97793	3.172	3.901 ₋₄	3.936	2.737 ₋₄	5.982	7.394 ₋₄	7.672	8.130 ₋₄
7	392192	3.172	1.024 ₋₄	3.936	6.855 ₋₅	5.982	1.949 ₋₄	7.671	2.035 ₋₄

TABLE 5. Data for the L-shape problem, corresponding to approximations of λ_1 through λ_4 , $\kappa = 0.2$ and $\kappa = 0.3$ for different singular points, $c = 1/2$.

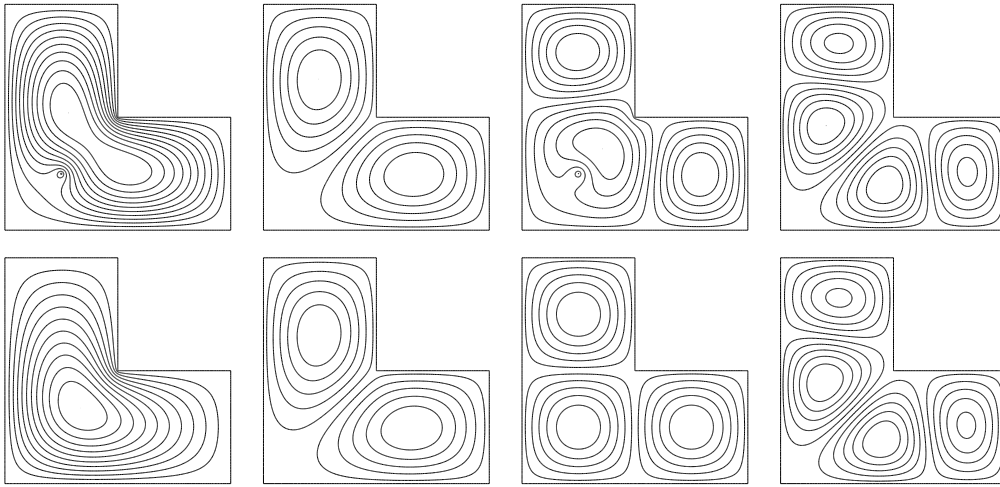


FIGURE 3. Contour plots of $\phi_{k,5}$ for $k = 1, 2, 3, 4$ for the L-shape problem. For the top row we have $c = 1/2$, and for the bottom row $c = 0$.

REFERENCES

- [1] T. Apel and S. Nicaise. The finite element method with anisotropic mesh grading for elliptic problems in domains with corners and edges. *Math. Methods Appl. Sci.*, 21(6):519–549, 1998.
- [2] I. Babuška, R. B. Kellogg, and J. Pitkäranta. Direct and inverse error estimates for finite elements with mesh refinements. *Numer. Math.*, 33(4):447–471, 1979.
- [3] I. Babuška and J. Osborn. Eigenvalue problems. In *Handbook of numerical analysis, Vol. II*, Handb. Numer. Anal., II, pages 641–787. North-Holland, Amsterdam, 1991.
- [4] C. Băcuță, V. Nistor, and L. T. Zikatanov. Improving the rate of convergence of ‘high order finite elements’ on polygons and domains with cusps. *Numer. Math.*, 100(2):165–184, 2005.
- [5] R. E. Bank. PLTMG: A software package for solving elliptic partial differential equations. Users’ Guide 10.0. Technical report, University of California at San Diego, Department of Mathematics, 2007.
- [6] R. E. Bank, L. Grubišić, and J. S. Owall. A framework for robust eigenvalue and eigenvector error estimation and ritz value convergence enhancement. *Applied Numer. Math.*, 66(0):1–29, 2013.
- [7] H. Chen, L. He, and A. Zhou. Finite element approximations of nonlinear eigenvalue problems in quantum physics. *Comput. Methods Appl. Mech. Engrg.*, 200(21-22):1846–1865, 2011.
- [8] M. Dauge. *Elliptic Boundary Value Problems on Corner Domains*, volume 1341 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 1988.
- [9] A. Ern and J.-L. Guermond. *Theory and practice of finite elements*, volume 159 of *Applied Mathematical Sciences*. Springer-Verlag, New York, 2004.
- [10] V. Felli, A. Ferrero, and S. Terracini. Asymptotic behavior of solutions to Schrödinger equations near an isolated singularity of the electromagnetic potential. *J. Eur. Math. Soc. (JEMS)*, 13(1):119–174, 2011.

- [11] V. Felli, E. Marchini, and S. Terracini. On the behavior of solutions to Schrödinger equations with dipole type potentials near the singularity. *Discrete Contin. Dyn. Syst.*, 21(1):91–119, 2008.
- [12] S. Fournais, M. Hoffmann-Ostenhof, T. Hoffmann-Ostenhof, and T. Østergaard Sørensen. Analytic structure of solutions to multiconfiguration equations. *J. Phys. A*, 42(31):315208, 11, 2009.
- [13] P. Grisvard. *Elliptic problems in nonsmooth domains*, volume 24 of *Monographs and Studies in Mathematics*. Pitman (Advanced Publishing Program), Boston, MA, 1985.
- [14] P. Grisvard. *Singularities in boundary value problems*, volume 22 of *Recherches en Mathématiques Appliquées [Research in Applied Mathematics]*. Masson, Paris, 1992.
- [15] E. Hunsicker, H. Li, V. Nistor, and U. Ville. Analysis of Schrödinger operators with inverse square potentials I: regularity results in 3D. *Bull. Math. Soc. Sci. Math. Roumanie (N.S.)*, 55(103)(2):157–178, 2012.
- [16] V. A. Kondrat'ev. Boundary value problems for elliptic equations in domains with conical or angular points. *Trudy Moskov. Mat. Obsč.*, 16:209–292, 1967.
- [17] V. A. Kozlov, V. G. Maz'ya, and J. Rossmann. *Elliptic boundary value problems in domains with point singularities*, volume 52 of *Mathematical Surveys and Monographs*. American Mathematical Society, Providence, RI, 1997.
- [18] V. A. Kozlov, V. G. Maz'ya, and J. Rossmann. *Spectral problems associated with corner singularities of solutions to elliptic equations*, volume 85 of *Mathematical Surveys and Monographs*. American Mathematical Society, Providence, RI, 2001.
- [19] R. B. Lehoucq, D. C. Sorensen, and C. Yang. *ARPACK users' guide*, volume 6 of *Software, Environments, and Tools*. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1998. Solution of large-scale eigenvalue problems with implicitly restarted Arnoldi methods.
- [20] H. Li, A. Mazzucato, and V. Nistor. Analysis of the finite element method for transmission/mixed boundary value problems on general polygonal domains. *Electron. Trans. Numer. Anal.*, 37:41–69, 2010.
- [21] H. Li and V. Nistor. Analysis of a modified Schrödinger operator in 2D: regularity, index, and FEM. *J. Comput. Appl. Math.*, 224(1):320–338, 2009.
- [22] H. Li and J. S. Ovall. A posteriori error estimation of hierarchical type for the schrödinger operator with inverse square potential. *In Review*, 2013.
- [23] S. Moroz and R. Schmidt. Nonrelativistic inverse square potential, scale anomaly, and complex extension. Preprint hep-th/0909.3477v3, 2010.
- [24] A. Naga and Z. Zhang. Function value recovery and its application in eigenvalue problems. *SIAM J. Numer. Anal.*, 50(1):272–286, 2012.
- [25] G. Strang and G. J. Fix. *An analysis of the finite element method*. Prentice-Hall Inc., Englewood Cliffs, N. J., 1973. Prentice-Hall Series in Automatic Computation.
- [26] L. N. Trefethen and T. Betcke. Computed eigenmodes of planar regions. In *Recent advances in differential equations and mathematical physics*, volume 412 of *Contemp. Math.*, pages 297–314. Amer. Math. Soc., Providence, RI, 2006.
- [27] N. M. Wigley. Asymptotic expansions at a corner of solutions of mixed boundary value problems. *J. Math. Mech.*, 13:549–576, 1964.
- [28] H. Wu and D. W. L. Sprung. Inverse-square potential and the quantum vortex. *Phys. Rev. A*, 49:4305–4311, Jun 1994.

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